

AD-A159 817

DEVELOPMENT OF SOLID STATE VIBRATIONAL LASERS(U)  
CORNELL UNIV ITHACA NY LAB OF ATOMIC AND SOLID STATE  
PHYSICS A J SIEVERS 01 AUG 85 ARO-21574 1-PH

1/1

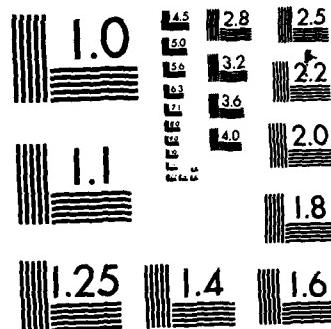
UNCLASSIFIED

DAAG29-84-K-0034

F/G 20/5

NL

END  
FILED  
DTIC



MICROCOPY RESOLUTION TEST CHART  
NATIONAL BUREAU OF STANDARDS-1963-A

AD-A159 817

UNCLASSIFIED  
SECURITY CLASSIFICATION OF THIS PAGE (When Data Entered)

**REPORT DOCUMENTATION PAGE**

1. REPORT NUMBER ARO 21574.1-PH		2. GOVT ACCESSION NO. N/A	3. RECIPIENT'S CATALOG NUMBER N/A
4. TITLE (and Subtitle) Development of Solid State Vibrational Lasers		5. TYPE OF REPORT & PERIOD COVERED Final Report 4/1/84 - 3/31/85	6. PERFORMING ORG. REPORT NUMBER DAAG29-84-K-0034
7. PERFORMING ORGANIZATION NAME AND ADDRESS Lab. of Atomic and Solid State Physics Cornell University Ithaca, NY 14853-2501		8. CONTRACT OR GRANT NUMBER(s) DAAG29-84-K-0034	9. PROGRAM ELEMENT, PROJECT, TASK AREA & WORK UNIT NUMBERS N/A
9. CONTROLLING OFFICE NAME AND ADDRESS U. S. Army Research Office Post Office Box 12211 Research Triangle Park NC 27709		10. REPORT DATE August 1, 1985	11. NUMBER OF PAGES 15
12. MONITORING AGENCY NAME & ADDRESS (if different from Controlling Office)		13. SECURITY CLASS. (of this report) Unclassified	14. DECLASSIFICATION/DOWNGRADING SCHEDULE SCHEDULE
15. DISTRIBUTION STATEMENT (of this Report) Approved for public release; distribution unlimited.			

DTIC  
ELECTED  
S OCT 08 1985  
D  
E

EXTRAMURAL DOCUMENTATION FORM  
1. EDITION OF 1 NOV 68 IS OBSOLETE  
2. SECURITY CLASSIFICATION  
UNCLASSIFIED  
3. DISTRIBUTION STATEMENT (of the abstract entered in Block 20, if different from Report)  
Approved for public release; distribution unlimited.

4. ELEMENARY NOTES  
view, opinions, and/or findings contained in this report are  
of the author(s) and should not be construed as an official  
ment of the Army position, policy, or decision, unless so  
nared by other documentation.  
SOS (Continue on reverse side if necessary and identify by block number)  
ious-wave lasing; vibrational transition; population inversion; optical  
; overtone; tunable; color center; fluorescence; vibrational energy exchange;  
onal laser; color center; fluorescence; vibrational energy exchange;  
transfer; inhomogeneous; Manley-Rowe  
SOS (Continue on reverse side if necessary and identify by block number)  
ious-wave lasing in the 5  $\mu\text{m}$  wavelength region has been obtained on  
brational transition of impurity CN<sup>-</sup> ions in KBr. Population inver-  
duced by optical pumping of the weakly allowed first-overtone level  
cule with a tunable (F<sub>2</sub>)<sup>+</sup>. A color center laser.  
investigated the energy budget of this solid-state vibrational  
level is pumped with a CW color center laser.

UNCLASSIFIED

SECURITY CLASSIFICATION OF THIS PAGE(When Data Entered)

20.

Using pulsed fluorescence techniques and numerical simulations, we find that due to vibrational energy exchange between pairs of CN<sub>2</sub> ions in the v = 1 lower laser level, "waste" pump energy is recycled to the v = 2 upper level. Ideally the efficiency of this recycling scheme approaches unity, but due to energy transfer into higher vibrational levels, and because of spectral diffusion within the inhomogeneously broadened CN<sub>2</sub> profile, the observed recycling efficiency is reduced. To our knowledge this is the first experiment demonstrating defeat of the Manley-Rowe limit. *Keywords:* -> to facilitate

UNCLASSIFIED

SECURITY CLASSIFICATION OF THIS PAGE(When Data Entered)

ARO 21574.1-PH

**Development of Solid State Vibrational Lasers**

4/1/84 - 3/31/85

**FINAL REPORT**

A. J. Sievers

Professor of Physics

Laboratory of Atomic and Solid State Physics

August 1, 1985

U. S. Army Research Office

DAAG29-83-K-0034

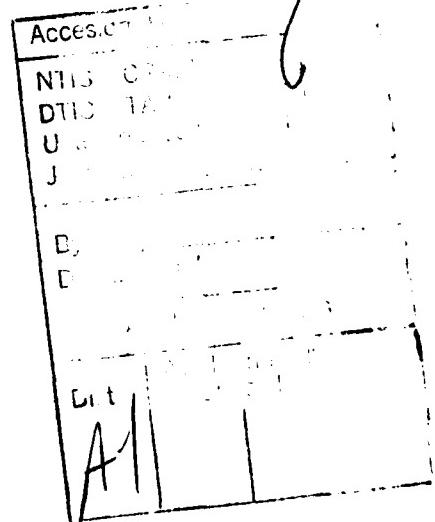
Cornell University

APPROVED FOR PUBLIC RELEASE:  
DISTRIBUTION UNLIMITED

85 10 08 031

## TABLE OF CONTENTS

	<u>Page</u>
I. Summary . . . . .	1
II. Statement of the Problem . . . . .	1
III. Experimental Investigation . . . . .	2
IV. Discussion and Conclusions . . . . .	4
V. References . . . . .	7
VI. Publications . . . . .	8
VII. Scientific Personnel . . . . .	9



## I. Summary

Continuous-wave lasing in the  $5 \mu\text{m}$  wavelength region has been obtained on the 2+1 vibrational transition of impurity CN<sup>-</sup> ions in KBr. Population inversion is produced by optical pumping of the weakly allowed first-overtone level of the molecule with a tunable (F<sub>2</sub><sup>+</sup>)<sub>A</sub> color center laser.

We have investigated the energy budget of this solid-state vibrational laser when the v = 2 upper laser level is pumped with a CW color center laser. Using pulsed fluorescence techniques and numerical simulations, we find that due to vibrational energy exchange between pairs of CN<sup>-</sup> ions in the v = 1 lower laser level, "waste" pump energy is recycled to the v = 2 upper level. Ideally the efficiency of this recycling scheme approaches unity, but due to energy transfer into higher vibrational levels, and because of spectral diffusion within the inhomogeneously broadened CN<sup>-</sup> profile, the observed recycling efficiency is reduced. To our knowledge this is the first experiment demonstrating defeat of the Manley-Rowe limit.

## II. Statement of the Problem

With the discovery of vibrational fluorescence from CN<sup>-</sup> ions embedded in KCl host matrices<sup>1</sup> came the first possibility for the construction of a solid-state laser based on purely vibrational transitions. This possibility was recently realized in KBr crystals doped with 0.5 mol % KCN when optical pumping between the v = 0 and v = 1 levels produced lasing between the v = 2 and v = 1 vibrational levels of CN<sup>-</sup> in a pulsed mode.<sup>2</sup> The crucial mechanism leading to a population inversion in this high concentration system is the vibrational energy exchange between proximate CN<sup>-</sup> ions in their v = 1 state. The result of this energy exchange (V-V transfer is the rapid depletion of the v = 1 lower laser level population in favor of the v = 2 upper laser level and the v = 0 ground state.

Although vibrational energy exchange between diatomic molecules in IR gas lasers has been well studied, the observation of laser action on a vibrational transition of CN<sup>-</sup> ions embedded in KBr<sup>2</sup> raises new questions about the role of V-V transfer processes in this different context. In addition reliance on this interesting and unusual pumping scheme imposes two important limitations on the laser: The V-V transfer process limits the device only to pulsed operation when pumped on the 0+1 transition, and second, at increased temperature V-V transfer depopulates the v = 2 state in favor of the v = 1 state, restricting laser operation to temperatures below 4 K.

In this final report we describe our successful attempt to overcome these limitations. A tunable (F<sub>2</sub><sup>+</sup>)<sub>A</sub> color-center laser is used to pump directly the weakly allowed 0+2 vibrational transition of the CN<sup>-</sup> ion at 2.4  $\mu\text{m}$  in KBr crystals doped with 0.5 and 0.05 mol % KCN. Continuous-wave lasing is observed to occur on the 2+1 transition at 4.9  $\mu\text{m}$  for both crystals.

### III. Experimental Investigation

The doped KBr crystals used in these measurements are prepared by the method described in Ref. 2. Resonant cavities are fabricated by evaporating reflective films directly onto the 1 mm x 1 mm end faces of cleaved crystals of different lengths. For 0.5% concentration a 5-mm long cavity is formed with a  $\lambda/4$  (at 4.9  $\mu\text{m}$ ) germanium layer on one end face and a gold layer on the other. Efficient input coupling of the 2.4- $\mu\text{m}$  pump light is possible through the germanium layer, which has an antireflection coating at the pump wavelength but as a 70% reflector at the lasing wavelength. For the 0.05% concentration, gold films are used on both ends of a 14-mm long cavity, with input coupling provided by a 0.2 mm hole in one of the films. For ease of optical alignment on both crystals, output coupling of a few percent is provided by a small hole in the gold end reflector opposite the input coupler.

The KBr crystal cavities are attached with wire to the bottom of a stage suspended in a liquid helium bath. The temperature of the laser cavities is variable down to 1.7 K. The color-center laser beam is focused by a 100 mm lens through two cryostat windows into the end mirror of the KBr cavity. No attempt was made to mode match the pump spot with the cavity, as the plane-parallel structure of the KBr lasers certainly permitted operation in numerous transverse modes. The output coupling hole on the opposite end of the laser is imaged either directly or through a 0.5 m monochromator onto a 100 kHz bandwidth InSb detector. Lasing is unmistakably detected by a sudden off-scale increase of the detected signal on a monitoring oscilloscope and by the rapid rise and fall times of the signal compared with the radiative lifetime (25 msec).

The color-center pump laser is based on the  $(F_2^+)_A$  center in lithium-doped KCl. The necessary pump wavelength ( $2.42 \mu\text{m}$ ) is near the extreme end of the tuning range of the color-center crystal, so output power is limited to a continuous-wave value of 13 mW. Laser construction is similar to that described in Ref. 3. Coarse tuning is achieved by placing a sapphire prism in the output arm of the color-center laser cavity and rotating the output coupler. This construction also separates residual Nd:YAG laser pump radiation from the color-center output, thus reducing the subsequent heat load on the liquid-helium cryostat. Fine tuning the color-center laser across the absorption profile of the  $0+2$  transition of the CN<sup>-</sup> ions is accomplished by angle tuning a 30% reflective 1 mm interacavity étalon. The final linewidth of this color center laser is less than 1 GHz, which corresponds to two or three active longitudinal modes of the cavity.

A level diagram depicting the overtone pumping scheme is shown in Fig.

1. From Fourier-transform spectrometer measurements the inhomogeneous line-width of the 0+2 transition is  $0.5 \text{ cm}^{-1}$ , and the absorption coefficient at line center is  $2.1 \text{ cm}^{-1}$  for the KBr sample containing 0.5 mol KCN at 1.7 K. The strength of this transition is roughly 1% that of the fundamental; however, it is still strong enough to absorb most of the pump energy in a 1-cm-long crystal. The gold mirrors on both ends of the lightly doped crystal help to trap the pump radiation, increasing the effective absorption.

Figure 2(a) shows the output from the 0.5 mol % cavity at 1.7 K (upper trace) along with the chopped output of the 13-mW pump laser. Note the brief delay following the leading edge of the pump pulse before oscillation begins. This delay is due to the necessary buildup of gain to reach threshold. The trailing edge of the output pulse closely follows the pump. This observation suggested the possibility of continuous-wave laser operation, and this was confirmed for both the 0.5 and 0.5% cavities. A typical trace is shown in Fig. 2(b). The output power for both cavities was similar and was estimated to be a few microwatts, although the output power through the germanium input coupler of the 0.5% cavity was probably several times larger.

One expected characteristic of the laser output for such a relatively long-lived upper state is relaxation oscillations. These are visible in the traces of the output power.

The temperature dependence of the gain in both cavities is found to be very similar to that of the pulse laser with no lasing occurring above 7 K.

#### **IV. Discussion and Conclusions**

The fact that cw operation of this device is observed is perhaps at first surprising, since a bottleneck is expected at the lower laser level due

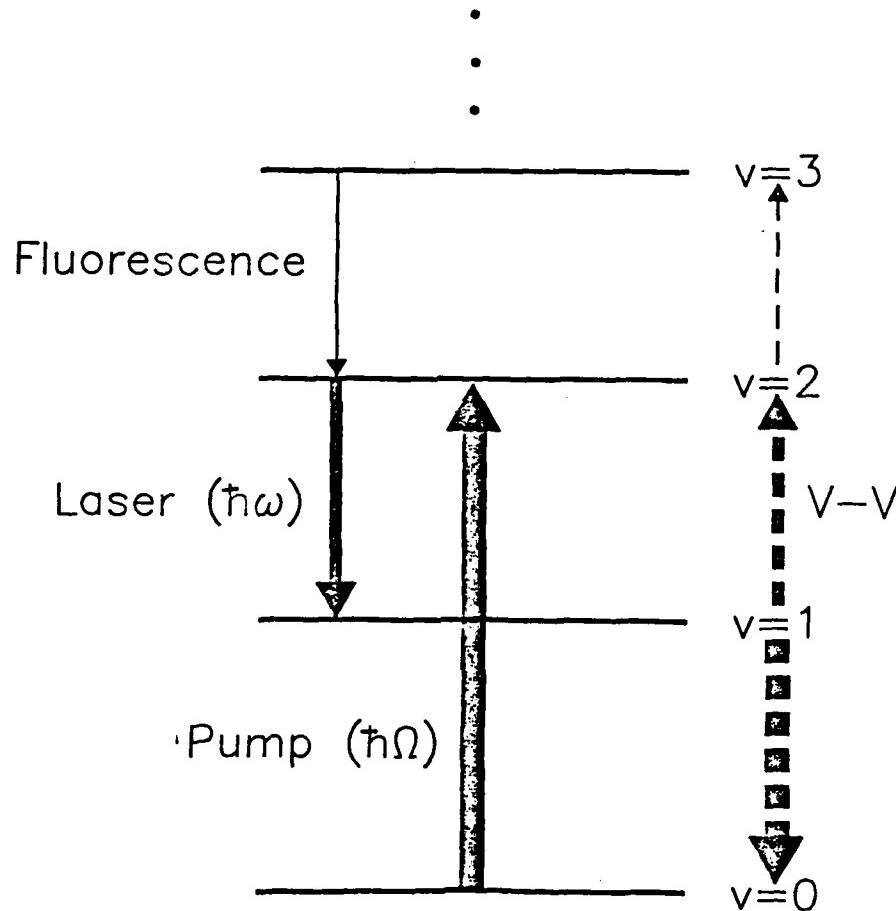


Figure 1: Energy level diagram of the anharmonic vibrational manifold of  $\text{CN}^-$  showing schematically the repopulating effects of energy exchange. Optical transitions are indicated as solid arrows, while dashed arrows show nonradiative transitions.

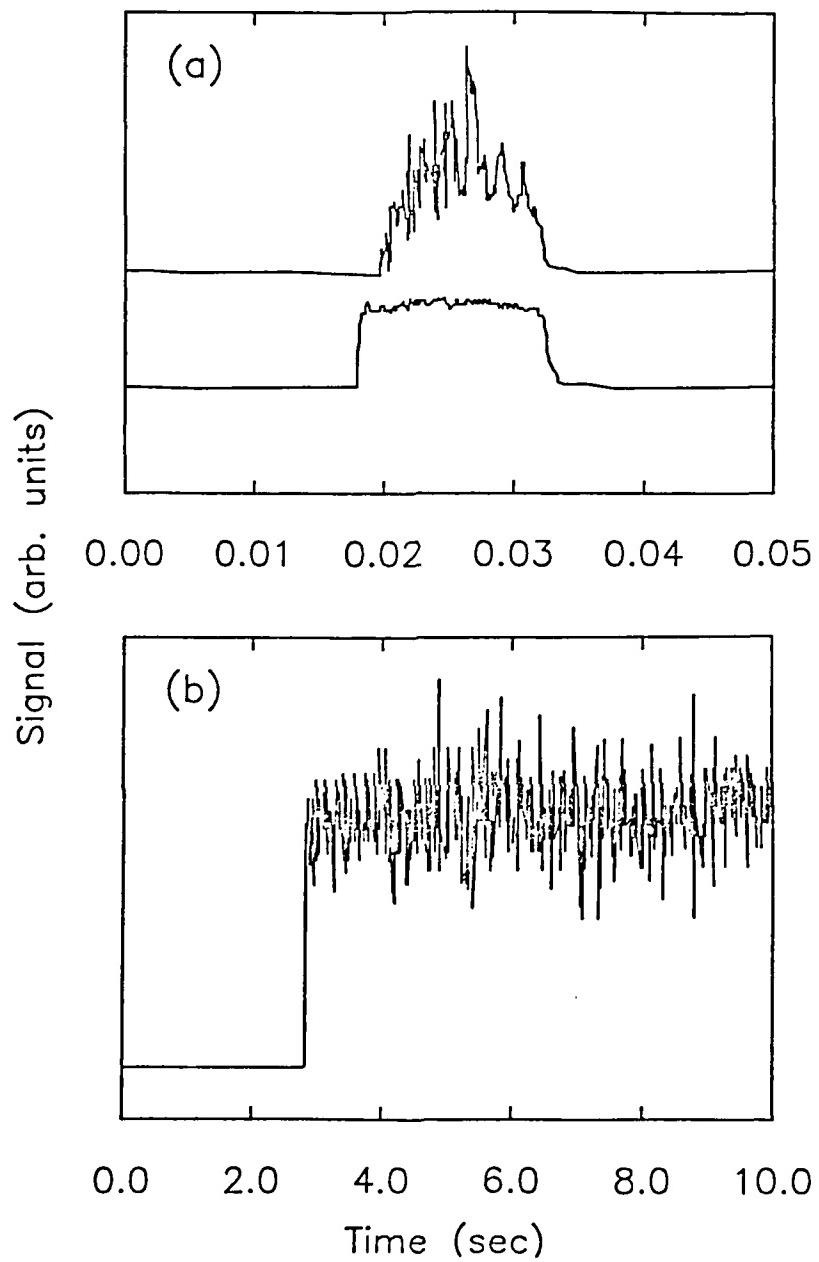


Figure 2: (a) Lasing at  $2054\text{ cm}^{-1}$  from a 0.5% KCN-doped crystal at 1.7 K (upper trace) versus the chopped pump pulse at  $4130\text{ cm}^{-1}$ . (b) Continuous-wave lasing of a 0.05% sample. Peak output power for both types of crystal is estimated at a few microwatts.

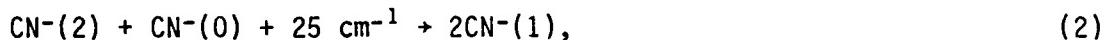
to the long radiative lifetime (~50 msec) of the  $v = 1$  state. However, rapid depletion of the  $v = 1$  state population by V-V transfer process (1) as well as by the general exchange reaction



reduces the effective lifetime of the lower laser level to <1 msec.

Direct evidence for the lifetime shortening effect of V-V transfer on the  $v = 1$  state is shown in Fig. 3. An initial population in the  $v = 1$  state is produced with a short pulse of doubled CO<sub>2</sub> laser radiation, and the subsequent fluorescence signals indicate the time evolution of the vibrational populations. Note the rapid fall in signal from the  $1 \rightarrow 0$  transition and the corresponding fast rise of the  $2 \rightarrow 1$  fluorescence signal.

The quenching of the laser gain is due to the thermally induced V-V transfer reaction<sup>2,4</sup>



which rapidly depletes the upper laser level. On the other hand, the spontaneous reverse process, which dominates below 4 K, acts to repopulate the  $n = 2$  level once the  $2 \rightarrow 1$  transition becomes saturated. Indeed, only weak fluorescence from the  $1 \rightarrow 0$  transition is observed at low temperatures. Below 4 K, this transfer back to the  $v = 2$  state reduces to less than 1 msec the effective lifetime of the lower laser level.<sup>5</sup>

An important consequence of overtone pumping is the following: Pump energy delivered to the  $v = 1$  level by gain saturation of the  $2+1$  transition is partly recycled to the upper laser level via process (2). A more complete

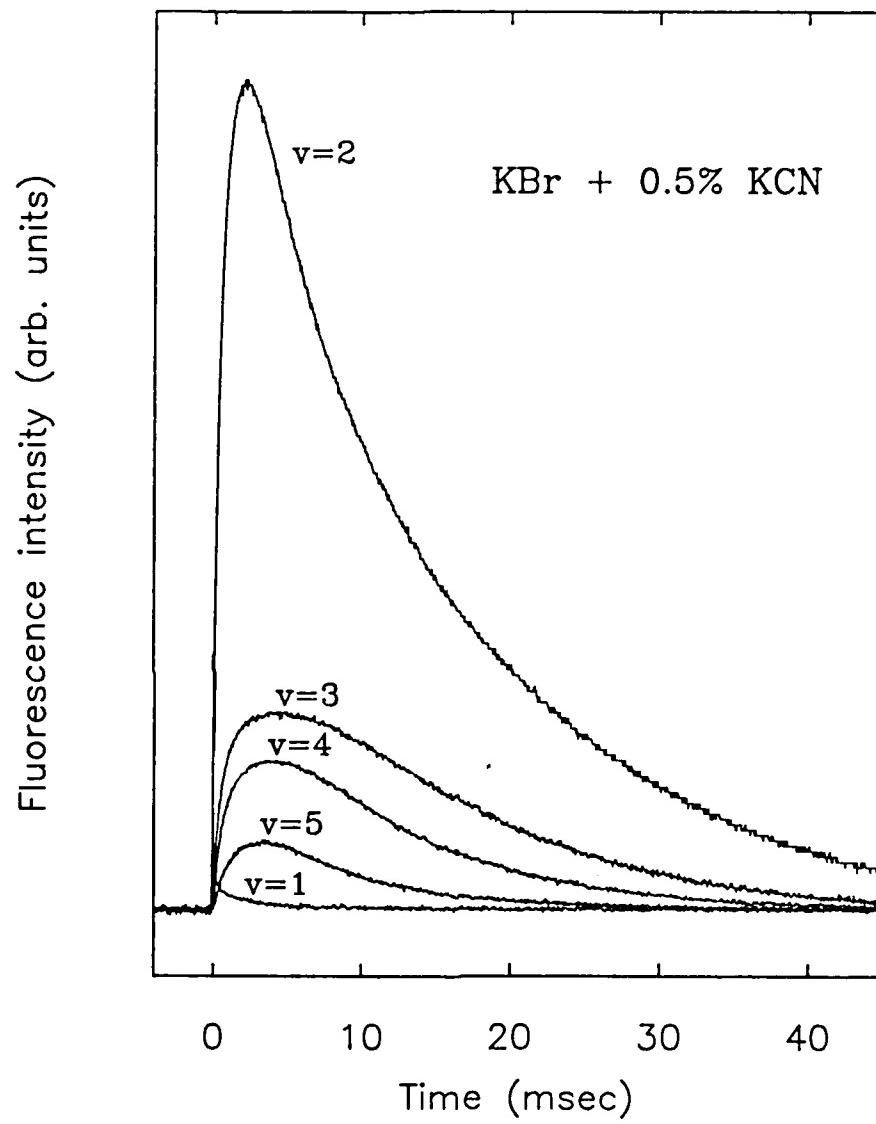


Figure 3: Fluorescence signals from the first five vibrational transitions of  $\text{CN}^-$  following pulsed pumping of the  $v = 1$  state. V-V transfer rapidly depletes the  $v = 1$  level.

accounting of the budget in these circumstances is illustrated in the level diagram of Fig. 1 which includes the influence of process (1). Normally, the efficiency of an optically pumped three-level laser is constrained by the Manley-Rowe limit  $\epsilon = \omega/\Omega$ , where  $\Omega$  is the pump frequency and  $\omega$  is the output frequency. In the KBr:CN<sup>-</sup> system, however, V-V transfer bypasses the Manley-Rowe limit, which in an ideal case would lead to efficiencies exceeding 50%. Although our observed overall efficiency is much less than this value due to external and internal effects unrelated to V-V transfer, the recycling process does play an important role. These other operating efficiencies aside, the KBr:CN<sup>-</sup> laser is the first experimental system to our knowledge to defeat the Manley-Rowe limit.

The results presented here demonstrate the feasibility of an alternative pumping scheme for the KBr:CN<sup>-</sup> laser that should lead to higher-temperature operation with broad tunability. Although the effect of thermal quenching of the gain through V-V transfer is still dominant at 0.05 mol % CN<sup>-</sup>, reducing the concentration even further will eventually decouple the CN<sup>-</sup> dipoles, permitting inversion above 4 K. Additionally, the rapidly increasing absorption linewidth of the ions resulting from the thermal occupation of hindered rotor levels<sup>6</sup> (FWHM of 50 cm<sup>-1</sup> at 50 K), combined with the measured high fluorescence quantum efficiency at 80 K (Ref. 7) and even at room temperature,<sup>8</sup> gives promising prospects for more convenient and useful operation of this device in the 5-μm spectral region.

V. References

1. Y. Yang and F. Lüty, Phys. Rev. Lett. 51, 419 (1983).
2. R. W. Tkach, T. R. Gosnell, and A. J. Sievers, Opt. Lett. 9, 122 (1984).
3. C. R. Pollock and D. A. Jennings, Appl. Phys. B 28, 308 (1982).
4. J. Manz, Chem. Phys. Lett. 51, 477 (1977).
5. T. R. Gosnell, R. W. Tkach, and A. J. Sievers, Infrared Phys. 25, 35 (1985).
6. W. D. Seward and V. Narayananuriti, Phys. Rev. 148, 463 (1966).
7. K. P. Koch, Y. Yang and F. Lüty, Phys. Rev. B 29, 5840 (1984).
8. T. R. Gosnell, R. W. Tkach, and A. J. Sievers, Solid State Commun. 53, 419 (1985).

**VI. Publications**

- a) "Spontaneous Emission and Laser Oscillation from Vibrational Levels of Molecular Defects in Solids", T. R. Gosnell, R. W. Tkach and A. J. Sievers, *Infrared Physics* 25, 35 (1985).
- b) "IR Lasing from Localized Vibrational Modes in the Defect Solid State", T. R. Gosnell, R. W. Tkach and A. J. Sievers, 9th International Conference on Infrared and Millimeter Waves Digest (Japan Society of Applied Physics, 1984), p. 141.
- c) "CW Operation of the KBr:CN<sup>-</sup> Solid State Vibrational Laser in the 5 μm Region", T. R. Gosnell, A. J. Sievers and C. R. Pollock, *Optics Letters* 10, 125 (1985).
- d) "Bypassing the Manley-Rowe Limit in the KBr:CN<sup>-</sup> Vibrational Laser by Energy Recycling", T. R. Gosnell, C. R. Pollock and A. J. Sievers, *Bulletin Amer. Phys. Soc.* 30, 410 (1985).
- e) "Influence of Energy Transfer on the CW Operation of the KBr:CN<sup>-</sup> Solid State Vibrational Laser", T. R. Gosnell, A. J. Sievers and C. R. Pollock, Conference on Lasers and Electro Optics (1985) p. 98.
- f) "IR Hole Burning and Laser Action for Molecular Defects in Ionic Crystals" A. J. Sievers, *Bull. Amer. Phys. Soc.* 30, 390 (1985).

VII. Scientific Personnel

R. Tkach, K. Blum, A. J. Sievers

**END**

**FILMED**

**11-85**

**DTIC**